

Foundation for Polish Science

INNOVATIVE ECONOMY



# 50th Karpacz Winter School of Theoretical Physics H<sub>2</sub> molecule as a two-site system with the *ab initio* optimization of single-particle wave functions in the correlated state revisited: Electron-proton coupling Andrzej P. Kądzielawa<sup>1†</sup>, Agata Bielas<sup>2</sup>, Marcello Acquarone<sup>3</sup>, Maciej Maśka<sup>2</sup>, Józef Spałek<sup>1</sup>

<sup>1</sup>Instytut Fizyki im. Mariana Smoluchowskiego, Uniwersytet Jagielloński, ul. Reymonta 4, PL-30059 Kraków, Poland <sup>2</sup>Instytut Fizyki, Uniwersytet Śląski, ul. Uniwersytecka 4, PL-40007 Katowice, Poland <sup>3</sup>Dipartimento di Fisica e Scienze della Terra dell'Università di Parma, I-43100 Parma, Italy

<sup>†</sup>kadzielawa@th.if.uj.edu.pl

#### **MOTIVATION**

We want to have a realistic model of the  $H_2$  molecule. We start from Hamiltonian (frame Hamiltonian) with all one- and two-body interactions.

The main aims are to:

- 1. solve exactly the two-site Hamiltonian (frame *Solu*tion),
- 2. calculate the expressions for the microscopic parameters from the first principles (frame *Microscopic Pa*rameters),

### HAMILTONIAN

 $\mathcal{H} = \epsilon \left( \hat{n}_1 + \hat{n}_2 \right) + t \sum \left( \hat{a}_{1\sigma}^{\dagger} \hat{a}_{2\sigma} + \hat{a}_{2\sigma}^{\dagger} \hat{a}_{1\sigma} \right)$  $+ U \left( \hat{n}_{1\uparrow} \hat{n}_{1\downarrow} + \hat{n}_{2\uparrow} \hat{n}_{2\downarrow} \right) - 2J \mathbf{S}_1 \mathbf{S}_2$  $+\left(K-\frac{J}{2}\right)\hat{n}_{1}\hat{n}_{2}+J\left(\hat{a}_{1\uparrow}^{\dagger}\hat{a}_{1\downarrow}^{\dagger}\hat{a}_{2\downarrow}\hat{a}_{2\uparrow}+h.c.\right)$  $+V\sum\left[\left(\hat{n}_{1\sigma}+\hat{n}_{2\sigma}\right)\left(\hat{a}_{1\bar{\sigma}}^{\dagger}\hat{a}_{2\bar{\sigma}}+\hat{a}_{2\bar{\sigma}}^{\dagger}\hat{a}_{1\bar{\sigma}}\right)\right],$ 

#### **PARAMETERS EVALUATION**

We choose the Slater 1s basis

$$\Psi_{i}\left(\mathbf{r}\right) = \sqrt{\frac{\alpha^{3}}{\pi}} e^{-\alpha|\mathbf{r}-\mathbf{R}_{i}|},$$

where  $\alpha$  is the inverse size of the orbital. To ensure orthogonality we use Wannier functions which in this case reduce the superposition of the atomic states, i.e.,

 $w_{i}\left(\mathbf{r}\right) = \beta \left(\Psi_{i}\left(\mathbf{r}\right) - \gamma \Psi_{j}\left(\mathbf{r}\right)\right),$ with the mixing parameters  $\beta = \frac{1}{\sqrt{2}} \sqrt{\frac{1+\sqrt{1-S^2}}{1-S^2}}$  and  $\gamma =$  $\frac{S}{1+\sqrt{1-S^2}}$ , where  $S = S(\alpha, R) \equiv \langle \Psi_1 | \Psi_2 \rangle$  is the atomic functions' overlap.



- 3. optimize the single-particle wave-function to minimize the energy,
- 4. calculate the electron–ion couplings (frame *Electron*-proton coupling),
- 5. evaluate the zero-point motion (frame *Results*).

 $\epsilon$  is a single-particle energy, t the hopping parameter, U the onsite Coulomb repulsion, K the inter-site Coulomb repulsion, Jthe exchange integral and *V* the so-called correlated hopping. We also use implicit form  $\mathcal{H} = \sum_i \Xi_i \hat{O}_i$ , where  $\Xi_i$  is symbolic notation of any of the microscopic parameters and  $\hat{O}_i$  the notation of the respective operator.

### **MICROSCOPIC PARAMETERS**

The microscopic parameters  $\epsilon = T_{11}$ ,  $t = T_{12}$ ,  $U = V_{1111}$ ,  $J = V_{1212}$ ,  $K = V_{1122}$  and  $V = V_{1112}$  correspond to oneand two-particle interactions

> $T_{ij} = \langle w_i | \mathcal{T} | w_j \rangle,$  $V_{ijkl} = \langle w_i w_j | \mathcal{V}_{12} | w_k w_l \rangle \,,$

where in atomic units  $\mathcal{T} = -\nabla^2 - 2/|\mathbf{r} - \mathbf{R}|$ , and  $\mathcal{V} = -\nabla^2 - 2/|\mathbf{r} - \mathbf{R}|$  $2/|{\bf r}-{\bf r}'|.$ 

**Figs:** Microscopic param- ping parameter is V < 0. eters  $\epsilon$ , t, U, K, J and V versus average interionic distance R. Note the convergence of the intersite Coulomb repul-



#### SOLUTION

Solutions of our electronic Hamiltonian are listed in the table below. For given interionic distance R, our procedure requires us to minimize the energy  $E_i = E_i^e + \frac{2}{R}$  with respect to  $\alpha$ .



i	Eigenvector $ i\rangle$	Eigenenergy $E_i^e$	
1	$\hat{a}_{1\uparrow}^{\dagger}\hat{a}_{2\uparrow}^{\dagger}\left 0\right\rangle$	$2\epsilon + K - J$	
2	$\begin{vmatrix} \hat{a}_{1\downarrow}^{\dagger} \hat{a}_{2\downarrow}^{\dagger}   0 \rangle$	$2\epsilon + K - J$	
3	$\frac{1}{\sqrt{2}} \left( \hat{a}_{1\uparrow}^{\dagger} \hat{a}_{2\downarrow}^{\dagger} + \hat{a}_{1\downarrow}^{\dagger} \hat{a}_{2\uparrow}^{\dagger} \right) \left  0 \right\rangle$	$2\epsilon + K - J$	
6	$\frac{1}{\sqrt{2}} \left( \hat{a}_{1\uparrow}^{\dagger} \hat{a}_{1\downarrow}^{\dagger} - \hat{a}_{2\uparrow}^{\dagger} \hat{a}_{2\downarrow}^{\dagger} \right) \left  0 \right\rangle$	$2\epsilon + U - J$	
+	$\frac{(K-U-\mathcal{D})+4 t+V  5\rangle}{\sqrt{2\mathcal{D}(\mathcal{D}-U+K)}}$	$2\epsilon + \frac{U+K}{2} + J + \frac{1}{2}\mathcal{D}$	
	$\frac{(U-K+\mathcal{D}) 4\rangle+4 t+V  5\rangle}{\sqrt{2\mathcal{D}(\mathcal{D}+U-K)}}$	$2\epsilon + \frac{U+K}{2} + J - \frac{1}{2}\mathcal{D}$	
$\mathcal{D} = \sqrt{(U - K)^2 + 16(t + V)^2},  4\rangle = \sqrt{2}^{-1} \left( \hat{a}_{1\uparrow}^{\dagger} \hat{a}_{2\downarrow}^{\dagger} - \hat{a}_{1\downarrow}^{\dagger} \hat{a}_{2\uparrow}^{\dagger} \right)  0\rangle,$			
and $ 5\rangle = \sqrt{2}^{-1} \left( \hat{a}_{1\uparrow}^{\dagger} \hat{a}_{1\downarrow}^{\dagger} + \hat{a}_{2\uparrow}^{\dagger} \hat{a}_{2\downarrow}^{\dagger} \right)  0\rangle.$			



### **ELECTRON-PROTON COUPLING**



#### **Figs:** Coupling constants $\xi_{\epsilon}$ , $\xi_t$ , $\xi_U$ , $\xi_K$ , $\xi_J$ and $\xi_V$ versus intersite distance *R*.

**Fig:** The energy  $E_i$  versus intersite distance R. Clearly, the spin-singlet state  $|-\rangle$  is the equilibrium state.

## **EVALUATION OF THE ZERO-POINT MOTION**

We introduce the zero-point motion by extending the Hamiltonian in terms of displacement  $\delta R$ 

 $\mathcal{H} \to \mathcal{H}_{\delta R} = \mathcal{H} + \delta \mathcal{H} + \mathcal{H}_{ZPM},$ 

where

$$\left\langle \delta \mathcal{H} \right\rangle = \sum_{i} \xi_{i} \delta R \left\langle \hat{O}_{i} \right\rangle - \frac{2}{R_{B}^{2}} \delta R,$$

where  $\xi_i = \delta \Xi_i / \delta R$  are the coupling coefficients (see frame *Electron*-proton coupling).

$$\langle \mathcal{H}_{ZPM} \rangle = 2 \frac{\delta p^2}{2M_{H^-}} + \sum_{i=2}^9 \frac{1}{i!} E_B^{(i)} \delta R^i + o\left(\delta R^{10}\right),$$

where  $\delta p$  is the momentum of the  $H^-$  ion with mass  $M_{H^-}$ . We evaluate the value of  $\delta p$  from the Heisenberg principle and we minimize  $\langle \mathcal{H}_{\delta R} \rangle$  with respect to  $\delta R$ .



Fig: Averages of the operator part in Hamiltonian, calculated in the ground-state versus distance *R*. They are of the order of unity.

#### ACKNOWLEDGMENTS



We find the binding energy and distance

 $R_B = 1.43017 a_0$ ,

 $E_B = -2.29587 Ry.$ 

We also evaluate the zero-point motion amplitude

 $|\delta \mathbf{R}| = 0.189028 \ a_0$ 

The work was realised in the Project TEAM awarded to our group by the Foundation for Polish Science (FNP) for the years 2011-2014. The partial support by the special Grant MAESTRO, No. DEC-2012/04/A/ST3/00342, from the National Science Center (NCN) for the years 2012-2017 should also be acknowledged. APK visit at the University of Parma has been supported by the grant of the Polish Ministry of Science and Higher Education, Grant No. 7150/E-338/M/2013.

#### REFERENCES

RESULTS

[1] J. Spałek, R. Podsiadły, W. Wójcik, and A. Rycerz, Phys. [2] A. P. Kądzielawa, J. Spałek, J. Kurzyk, and W. Wójcik, [3] M. Acquarone and C. Noce, Int. J. of Mod. Phys. B 13, Rev. B **61**, 15676 (2000). Eur. Phys. J. B 86, 252 (2013). 3331 (1999).