

# $(H_2)_n$ molecule system with an *ab initio* optimization of wave functions in correlated state: Electron-proton couplings and intermolecular microscopic parameters

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# Outline

## 1 Objective

## 2 Methods

- Single-particle basis optimization

- Hamiltonian

- Statistically-consistent Gutzwiller Approximation (SGA)

## 3 Results

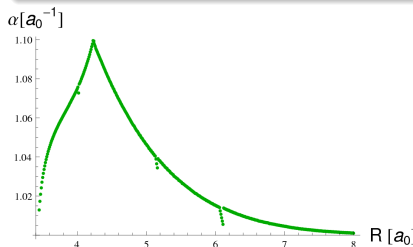
- Metal-insulator transition

- Quantum critical behavior

- Metallization pressure

## Objective

To apply SGA method for wave function quantum scaling problem near Mott-Hubbard transition.



J. Spałek, J. Kurzyk, R. Podsiadły,  
W. Wójcik, Eur. Phys. J. B **74**, 63-74  
(2010)

## Also

- to include the effect of external magnetic field on metal-insulator transition
- to calculate metallization pressure for atomic hydrogen

## Base functions

We build our basis by orthogonalizing Slater-type orbitals  $\Psi_i$  by introducing mixing coefficients  $\beta_j$

$$\langle w_{\pi_i(j)}(\mathbf{r}) | w_i(\mathbf{r}) \rangle = \delta_{i\pi_i(j)}, \quad (1)$$

$$w_i(\mathbf{r} - \mathbf{R}_i) = w_i(\mathbf{r}) = \sum_{j=0}^Z \beta_j \Psi_{\pi_i(j)}(\mathbf{r}), \quad (2)$$

where  $\pi_i$  is the  $Z$ -neighbors mapping function for the node  $i$ , where  $\pi_i(0) = i$ .

$$\Psi_i(\mathbf{r}) = \sqrt{\frac{\alpha^3}{\pi}} e^{-\alpha|\mathbf{r}-\mathbf{R}_i|}. \approx \alpha^{\frac{3}{2}} \sum_{a=1}^p B_a \left( \frac{2\Gamma_a^2}{\pi} \right)^{\frac{3}{4}} e^{-\Gamma_a^2|\mathbf{r}-\mathbf{R}_i|^2}. \quad (3)$$

# Microscopic parameters

## Values

We have two classes of microscopic parameters

$$T_{ij} = \langle w_i | \mathcal{H}_1 | w_j \rangle, \quad (4a)$$

$$V_{ijkl} = \left\langle w_i w_j \left| \frac{2}{|\mathbf{r} - \mathbf{r}'|} \right| w_k w_l \right\rangle, \quad (4b)$$

As we choose  $w_i$  to be real all possibilities can be described by

$\epsilon_i$	$t_{ij}$	$U_i$	$J_{ij}$	$K_{ij}$	$V_{ij}$
$T_{ii}$	$T_{ij}$	$V_{iiii}$	$V_{ijjj}$	$V_{ijij}$	$V_{iijj}$

# Hamiltonian

We start from the extended Hubbard model:

$$\begin{aligned}
 \mathcal{H} = & \sum_i \epsilon_i \hat{n}_i + \sum_{ij\sigma} t_{ij} \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \sum_i U_i \hat{n}_{i\uparrow} \hat{n}_{i\downarrow} - \sum_{ij} J_{ij} \mathbf{S}_i \mathbf{S}_j \quad (5) \\
 & + \frac{1}{2} \sum_{ij} \left( K_{ij} - \frac{J_{ij}}{2} \right) \hat{n}_i \hat{n}_j + \sum_{ij} J_{ij} \hat{c}_{i\uparrow}^\dagger \hat{c}_{i\downarrow}^\dagger \hat{c}_{j\downarrow} \hat{c}_{j\uparrow} \\
 & + \sum_{ij\sigma} V_{ij} \hat{n}_{i\sigma} \left( \hat{c}_{i\bar{\sigma}}^\dagger \hat{c}_{j\bar{\sigma}} + \hat{c}_{j\bar{\sigma}}^\dagger \hat{c}_{i\bar{\sigma}} \right),
 \end{aligned}$$

where  $\hat{c}_{i\sigma}$  and  $\hat{c}_{i\sigma}^\dagger$  are the fermionic operators of annihilation and creation of the electron with spin  $\sigma$  on 1s orbital of hydrogen atom  $i \bmod 2$  in  $H_2$  molecule  $[\frac{i}{2}]$ .

## SGA:

## Ground-state energy

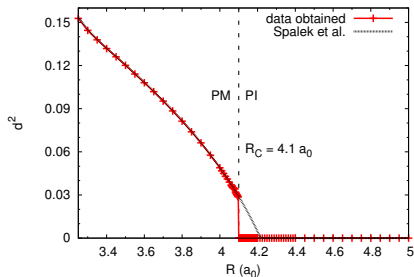
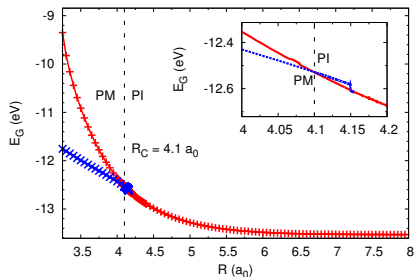
$$\frac{E_G}{\Lambda} = \epsilon_a^{eff} + \frac{1}{\Lambda} \left( \sum_{ij\sigma} t_{ij} \langle a_{i\sigma}^\dagger a_{j\sigma} \rangle + \sum_i U \langle a_{i\uparrow}^\dagger a_{i\uparrow} a_{i\downarrow}^\dagger a_{i\downarrow} \rangle \right) \quad (6)$$

⇓ minimization

$$\begin{aligned} \mathcal{K} = & \epsilon_a^{eff} \sum_{i\sigma} n_{i\sigma} + \sum_{ij\sigma} t_{ij} q_\sigma a_{i\sigma}^\dagger a_{j\sigma} + \Lambda U d^2 - \mu \sum_{i\sigma} n_{i\sigma} \\ & - \lambda_m \sum_i (m_i - m) - \lambda_n \sum_i (n_i - n) \end{aligned} \quad (7)$$

$n$	$m$	$d^2$	$\lambda_n$	$\lambda_m$	$\mu$
band filling	magnetization	no. of double occupancies	molecular field coupled with $n$	molecular field coupled with $m$	chemical potential

# Nature of transition



**Figure: G:** Ground-state energy of metallic (blue) and insulating (red) state. **D:** Double occupancies per site for our model vs. last publication.

1<sup>st</sup> order

Discontinuity of first derivative of energy (upper plot)  $\rightarrow$  metal-insulator transition of weakly first order.



## Zero-point motion for ion lattice

The uncertainties of momentum  $\delta P$  and distance between ions  $\delta R$ :

$$\Delta E = \frac{(\delta P)^2}{2M} + \sum_{i=1}^3 \frac{1}{2} \left( \frac{1}{R + \delta R^i} + \frac{1}{R - \delta R^i} \right) \quad (8)$$

using Heisenberg Principle  $(\delta \mathbf{P})(\delta \mathbf{R}) \geq \frac{3}{4}$  and minimizing with respect to  $\delta \mathbf{R}$ , we obtain results (at MIT):

	axis-aligned	plane-aligned	<b>diagonal</b>
$ \delta \mathbf{R}  (a_0)$	0.3432	0.3438	<b>0.3440</b>
$\Delta E - \frac{3}{R} (Ry)$	0.003455	0.003449	<b>0.003447</b>

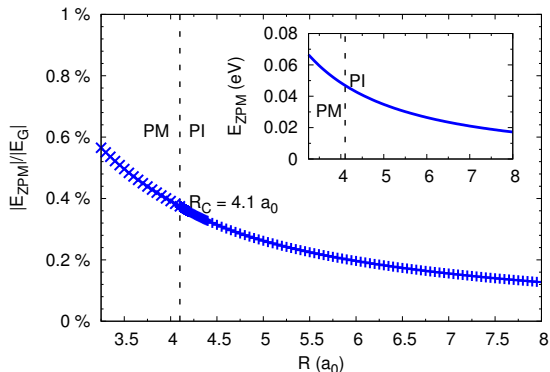
### Remark 1

All the calculations are done in atomic units.

### Remark 2

This is only estimation of the magnitude of ZPM.

# Zero-point motion



**Figure:** The relative magnitude of the zero-point motion energy with respect to ground state vs. lattice parameter  $R$ .

**Inset:** The explicit value of the energy of zero-point motion vs. lattice parameter  $R$ .

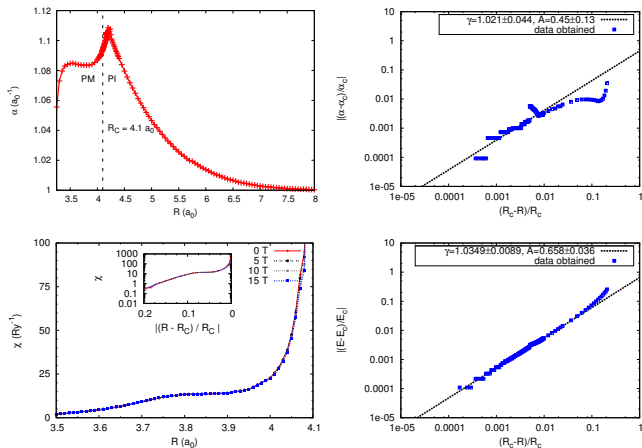


Figure:

**left:** reverse wave function size  $\alpha$  (**top**), and magnetic susceptibility (**bottom**) vs. lattice parameter  $R$ ,

**right:** scaling of reverse wave function size  $\alpha$  (**top**) and energy  $E$  (**bottom**) near critical point ( $\sim A((R_C - R)/R_C)^\gamma$ ).

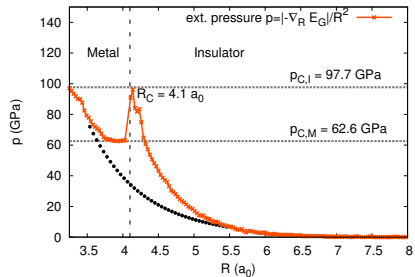
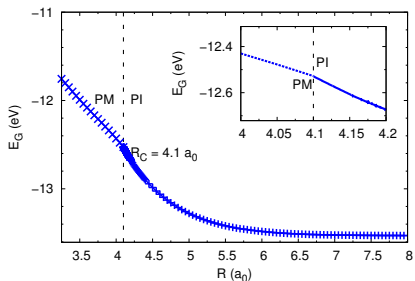


Figure: **top:** Ground energy vs. lattice parameter  $R$ , **bottom:** crystal stabilizing pressure vs. lattice parameter  $R$  - numerical (orange) and naive fit (black).

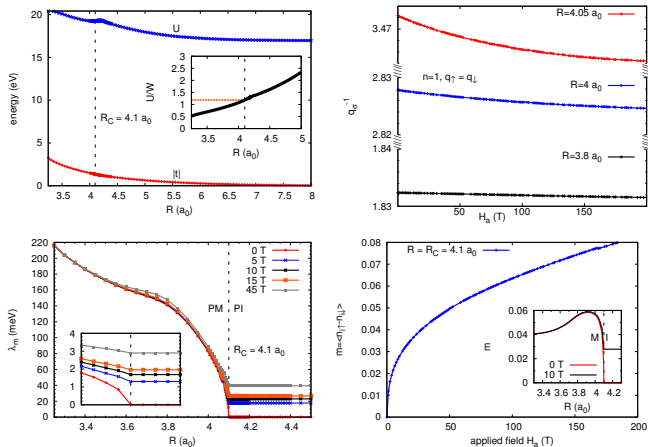
Pressure stabilizing hydrogen crystal

$$p_C = 97.7 \text{ GPa}$$

APK, J. Spałek, K. Kurzyk, W. Wójcik, Eur. Phys. J. B **86**, 252 (2013)

# Thank you!





**Figure:** **TL:** microscopic parameters vs. lattice parameter  $R$ , **TR:** effective mass enhancement vs. external magnetic field  $H_a$ , **BL:** effective magnetic field  $\lambda_m$  vs. lattice parameter  $R$ , **BR:** magnetization  $m$  vs. external magnetic field  $H_a$

**Table:** Values calculated using SGA method as a function of lattice parameter for SC. Units, if not written explicitly, are set to be *rydbergs* ( $Ry$ ).  $\chi(0)$  for  $R \geq R_c$  is infinite.

$R(a_0)$	$E_G^{SGA}$	$E_G^{GA}$	$t$	$U$	$\alpha^{-1}(a_0)$	$d^2$	$\lambda_m$	$\chi(Ry^{-1})$	$q^{-1}$
3.25	-0.8640	-0.8644	-0.2409	1.4996	0.9474	0.152774	0.015884	0.1809	1.17728
3.50	-0.8814	-0.8816	-0.1773	1.4749	0.9220	0.120128	0.012641	2.0598	1.36818
4.00	-0.9136	-0.9136	-0.1098	1.4152	0.9200	0.048886	0.006084	22.6577	2.82781
4.05	-0.9171	-0.9171	-0.1046	1.4139	0.9175	0.038973	0.004256	47.3562	3.47235
4.09	-0.9200		-0.1005	1.4140	0.9147	0.030193	0.027281	253.7567	4.40375
4.10	-0.9209	-0.9207	-0.0995	1.4143	0.9138	0.000000	0.000000	$\infty$	$\infty$
4.20	-0.9315	-0.9288	-0.0896	1.4217	0.9021	0.000000	0.000000		
4.50	-0.9544	-0.9517	-0.0705	1.3742	0.9263	0.000000	0.000000		
5.00	-0.9760	-0.9732	-0.0471	1.3200	0.9556	0.000000	0.000000		
7.00	-0.9939		-0.0082	1.2504	0.9972	0.000000	0.000000		
$\infty$	-1.0000	-1.0000	0.0000	1.2500	1.0000	0.000000	0.000000		

Minimization of free energy functional  $\mathcal{F}$ :

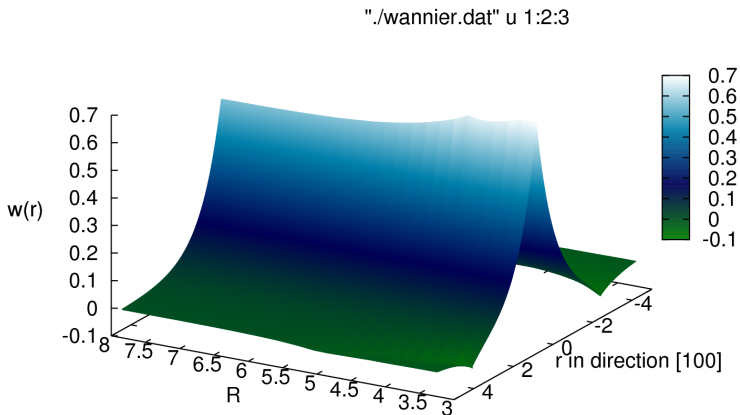
$$\mathcal{F}^{(SGA)} = -\frac{1}{\beta} \sum_{\mathbf{k}\sigma} \log \left( 1 + e^{-\beta E_{\mathbf{k}\sigma}^{(SGA)}} \right) + \Lambda (\lambda_n n + \lambda_m m + U d^2 + \mu n), \quad (9)$$

where the test eigenvalues  $E_{\mathbf{k}\sigma}^{(SGA)}$

$$E_{\mathbf{k}\sigma}^{(SGA)} \equiv q_\sigma \varepsilon_{\mathbf{k}} - \sigma (h + \lambda_m) - (\mu + \lambda_n),$$

$$q_\sigma = \frac{\left( \sqrt{(n_\sigma - d^2)(1 - n_\sigma - n_{\bar{\sigma}} + d^2)} + d \sqrt{n_{\bar{\sigma}} - d^2} \right)^2}{n_\sigma (1 - n_\sigma)} \quad (10)$$





**Figure:** Overall space profiles of the renormalized Wannier function for sc lattice as a function of lattice parameter  $R$  and along [100] direction  $r$ .