

Ab initio methods in Mott transition - application to metallization of atomic hydrogen

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1 Aims

2 Methods

- Hamiltonian
- Single-particle basis optimization
- Statistically-consistent Gutzwiller Approximation (SGA)

3 Results

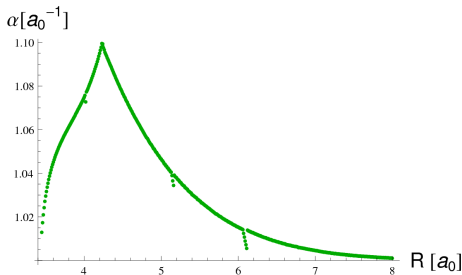
- Transition
- Microscopic parameters
- Results properties: quantum critical behavior
- Metallization pressure

4 Future

Aim 1

To apply SGA method for wave function quantum scaling problem near Mott-Hubbard transition.

For GA: Eur. Phys. J. B **74**, 63-74 (2010)



Other

- include the effect of external magnetic field on metal-insulator transition
- calculate metallization pressure for atomic hydrogen

We start from the Extended Hubbard Model:

$$\mathcal{H}_{EH} = \epsilon_a^{eff} \sum_i n_i + \sum_{i \neq j, \sigma} t_{ij} a_{i\sigma}^\dagger a_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow} + \frac{1}{2} \sum_{i \neq j} K_{ij} \delta n_i \delta n_j, \quad (1)$$

where effective energy

$$\epsilon_a^{eff} = \epsilon_a + \frac{1}{2\Lambda} \sum_{i \neq j} \left(K_{ij} + \frac{2}{|R_j - R_i|} \right), \quad (2)$$

t_{ij} – hopping integral, U – Coulomb interaction on site, a K_{ij} Coulomb interaction between site

External field H_a

Zeeman term:

$$\mathcal{H}_Z = - \sum_{i,\sigma} \sigma h n_{i\sigma}, \quad (3)$$

where $h = \frac{1}{2}g\mu_B H_a$ is reduced magnetic field.

Single-particle basis optimization

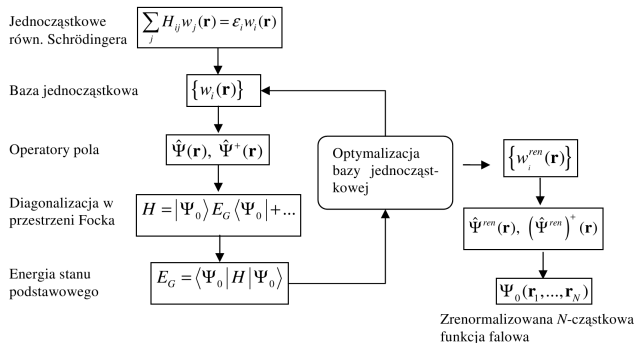


Figure: Flow chart from Dr. J. Kurzyk PhD thesis (UJ 2010)

On site $i = R_i$:

$$w_i(\vec{r}) = \beta \Psi_i(\vec{r}) - \gamma \sum_{j=1}^z \Psi_j(\vec{r}),$$

$$\Psi_i(\vec{r}) = \sqrt{\frac{\alpha^3}{\pi}} e^{-\alpha|\vec{r}-\vec{R}_i|} \quad (4)$$

$$\approx \alpha^{\frac{3}{2}} \sum_{a=1}^p B_a \left(\frac{2\Gamma_a^2}{\pi} \right)^{\frac{3}{4}} e^{-\Gamma_a^2|\vec{r}-\vec{R}_i|^2}.$$

constant	overlap	minimization $\mathcal{H}_1 \stackrel{a.u.}{=} -\nabla^2 - 2 \vec{r}-\vec{R}_i ^{-1}$	minimization
	dependent		with respect to E_G
z, p	β, γ	B_a, Γ_a	α

Values

Microscopic parameters ϵ_a , t_{ij} , U , and K_{ij} are expressed by integrals

$$\begin{aligned}t_{ij} &= \langle w_i | \mathcal{H}_1 | w_j \rangle, \\U &= \left\langle w_i w_i \left| \frac{e^2}{|\vec{r}_1 - \vec{r}_2|} \right| w_i w_i \right\rangle, \\K_{ij} &= \left\langle w_i w_j \left| \frac{e^2}{|\vec{r}_1 - \vec{r}_2|} \right| w_i w_j \right\rangle, \\ \epsilon_a &= \langle w_i | \mathcal{H}_1 | w_i \rangle.\end{aligned}\tag{5}$$

Minimization of free energy functional \mathcal{F} :

$$\mathcal{F}^{(SGA)} = -\frac{1}{\beta} \sum_{\vec{k}\sigma} \log \left(1 + e^{-\beta E_{\vec{k}\sigma}^{(SGA)}} \right) + \Lambda (\lambda_n n + \lambda_m m + U d^2 + \mu n), \quad (6)$$

where test eigenvalues $E_{\vec{k}\sigma}^{(SGA)}$

$$E_{\vec{k}\sigma}^{(SGA)} \equiv q_\sigma \varepsilon_{\vec{k}} - \sigma (h + \lambda_m) - (\mu + \lambda_n),$$

$$q_\sigma = \frac{\left(\sqrt{(n_\sigma - d^2)(1 - n_\sigma - n_{\bar{\sigma}} + d^2)} + d \sqrt{n_{\bar{\sigma}} - d^2} \right)^2}{n_\sigma (1 - n_\sigma)} \quad (7)$$

β	n	m	d	λ_n	λ_m	μ
$\frac{1}{k_B T}$	occup. density	magne-tization	num. doub. occupancy	mol. field con. w. n	mol. field con. w. m	chem. pot.

Ground energy

$$\frac{E_G}{\Lambda} = \epsilon_a^{\text{eff}} + \frac{1}{\Lambda} \left(\sum_{ij\sigma} t_{ij} \langle a_{i\sigma}^\dagger a_{j\sigma} \rangle + \sum_i U \langle a_{i\uparrow}^\dagger a_{i\uparrow} a_{i\downarrow}^\dagger a_{i\downarrow} \rangle \right) \quad (8)$$

↓ minimization

$$\begin{aligned} \mathcal{K} = \epsilon_a^{\text{eff}} \sum_{i\sigma} n_{i\sigma} + \sum_{ij\sigma} t_{ij} q_\sigma a_{i\sigma}^\dagger a_{j\sigma} + \Lambda U d^2 - \mu \sum_{i\sigma} n_{i\sigma} \\ - \lambda_m \sum_i (m_i - m) - \lambda_n \sum_i (n_i - n) \end{aligned} \quad (9)$$

Remark

No Zero Point Motion of protons.

Type of transition

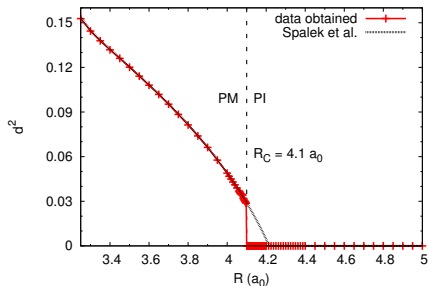
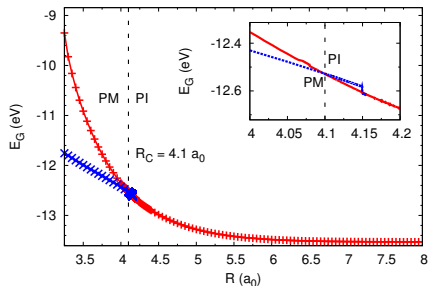


Figure: **G:** Ground energy of metallic (blue) and insulating (red) state. **D:** Double occupancies per site for our model vs. last publication.

I order

Discontinuity of first derivative of energy (upper plot) \rightarrow metal – insulator transition of weakly first order.

Zero Point Motion for ion chain

The uncertainties of momentum δP and distance between ions δR :

$$\Delta E = \frac{(\delta P)^2}{2M} + \frac{1}{2} \left(\frac{e^2}{R + \delta R} + \frac{e^2}{R - \delta R} \right) \stackrel{\delta R \ll R}{\approx} \frac{(\delta P)^2}{2M} + \frac{e^2}{R} + 2 \frac{e^2}{R} \left(\frac{\delta R}{R} \right)^2 \quad (10)$$

using Heisenberg Principle $(\delta P)(\delta R) \geq \frac{3}{4} \hbar^2$ and minimizing with respect to δR

$$\delta R = \left(\frac{3\hbar^2 R^3}{8Me^2} \right)^{\frac{1}{4}} \quad (11)$$

$$\Delta E = \frac{e^2}{R} + \hbar \left(\frac{6e^2}{MR^3} \right)^{\frac{1}{2}}$$

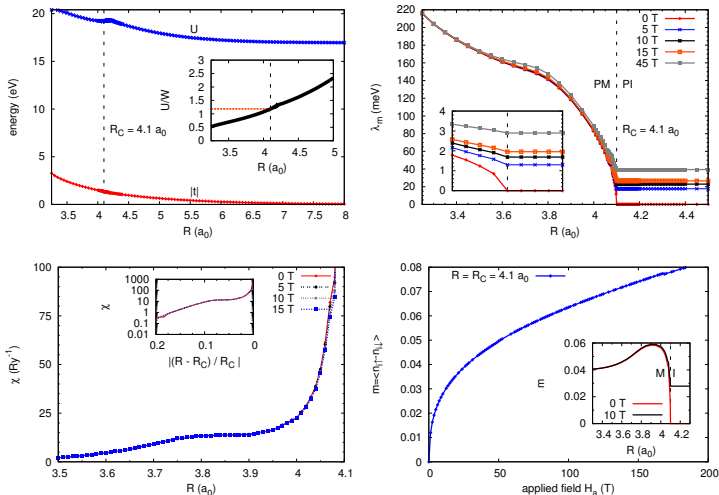


Figure: **GL:** microscopic parameters vs. lattice parameter R , **GP:** effective magnetic field λ_m vs. lattice parameter R , **DL:** magnetic susceptibility vs. lattice parameter R , **DP:** magnetization m vs. external magnetic field H_a

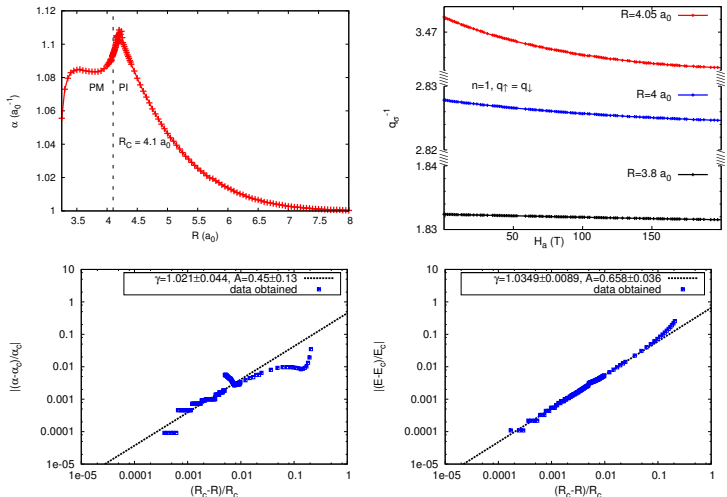


Figure: **GL:** reverse wave function size α vs. lattice parameter R , **GP:** effective mass enhancement vs. external magnetic field H_a , **D:** scaling of reverse wave function size $\alpha(L)$ and energy $E(P)$ near critical point.

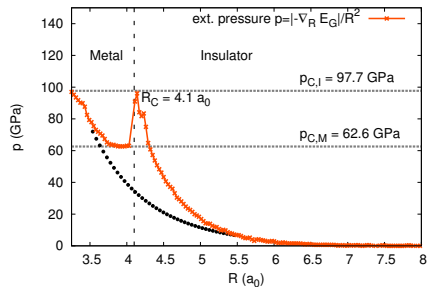
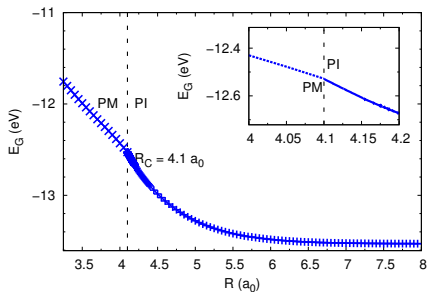


Figure: **G**: Ground energy vs. lattice parameter R , **D**: crystal stabilizing pressure vs. lattice parameter R - numerical (orange), naive approach with no phase transition (czarny).

Pressure stabilizing hydrogen crystal

$$p_C = 97.7 \text{ GPa}$$

Remark

To do: bulk modulus

$$\kappa = -\frac{1}{V} \frac{dV}{dp}$$

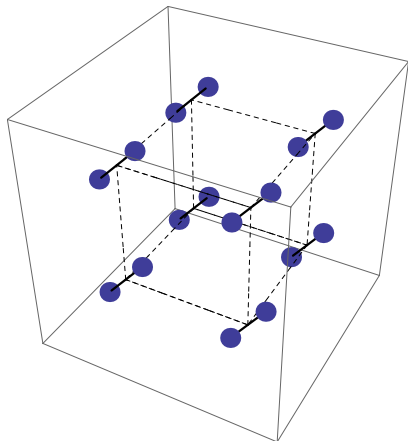
Modification 1

More realistic crystal structures:

- fcc
- hcp (antiferromagnetism ?)

Modification 2

Metal-insulator transition for
molecular hydrogen



Phononic part

$$\mathcal{H}_{ph} = \hbar\Omega_d \left(d^\dagger d + \frac{1}{2} \right) + \hbar\Omega_e \left(e^\dagger e + \frac{1}{2} \right) \quad (12)$$

Holstein term

$$\mathcal{H}_H = \hbar\Omega_d \xi \sum_{\sigma} \left(c_{1\sigma}^\dagger c_{1\sigma} - c_{2\sigma}^\dagger c_{2\sigma} \right) \left(d^\dagger + d \right) \quad (13)$$

Su-Schrieffer-Heeger term

$$\mathcal{H}_{SSH} = \hbar\Omega_e \zeta \sum_{\sigma} \left(c_{1\sigma}^\dagger c_{2\sigma} + c_{2\sigma}^\dagger c_{1\sigma} \right) \left(e^\dagger + e \right) \quad (14)$$



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Table: Values calculated using SGA method as a function of lattice parameter for SC. Units, if not written explicitly, are set to be *rydbergs* (Ry). $\chi(0)$ for $R \geq R_c$ is infinite.

$R(a_0)$	E_G^{SGA}	E_G^{GA}	t	U	$\alpha^{-1}(a_0)$	d^2	λ_m	$\chi(Ry^{-1})$	q^{-1}
3.25	-0.8640	-0.8644	-0.2409	1.4996	0.9474	0.152774	0.015884	0.1809	1.17728
3.50	-0.8814	-0.8816	-0.1773	1.4749	0.9220	0.120128	0.012641	2.0598	1.36818
4.00	-0.9136	-0.9136	-0.1098	1.4152	0.9200	0.048886	0.006084	22.6577	2.82781
4.05	-0.9171	-0.9171	-0.1046	1.4139	0.9175	0.038973	0.004256	47.3562	3.47235
4.09	-0.9200		-0.1005	1.4140	0.9147	0.030193	0.027281	253.7567	4.40375
4.10	-0.9209	-0.9207	-0.0995	1.4143	0.9138	0.000000	0.000000	∞	∞
4.20	-0.9315	-0.9288	-0.0896	1.4217	0.9021	0.000000	0.000000		
4.50	-0.9544	-0.9517	-0.0705	1.3742	0.9263	0.000000	0.000000		
5.00	-0.9760	-0.9732	-0.0471	1.3200	0.9556	0.000000	0.000000		
7.00	-0.9939		-0.0082	1.2504	0.9972	0.000000	0.000000		
∞	-1.0000	-1.0000	0.0000	1.2500	1.0000	0.000000	0.000000		